

Weak Segregation in Molten Statistical Copolymers

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Summary : we present a qualitative argument suggesting that a statistical (AB) copolymer should display a certain form of weak segregation, with A rich and B rich regions of size comparable to the coil radius. This should occur if the Flory parameter χ is larger than *unity* (much larger than for the phase separation of A and B of homopolymers) in agreement with detailed calculations by Frederikson, Milner and Leibler. We achieve a certain qualitative insight for the resulting microphase.

Keywords : microphase separation; polymer melts; segregation; statistical copolymers;

Introduction

Two distinct homopolymers (A...A) and (B...B) are usually non miscible, when the product of the Flory parameter χ by the polymerisation index N is large ($\chi N > 2$).^[1-3] This may be understood qualitatively through the work of transfer W of one A chain from a pure A medium to a pure B medium. This work is proportional to the chain length $W = NU$, where U is the enthalpy increase for the transfer of one A monomer ($U = \chi kT$, where T is the temperature). Segregation will occur when $W \gg kT$, or $\chi N \gg 1$.

In this note, we consider the more complex problem of a melt based on a statistical linear polymer AB AA B... with N monomers, and a completely random chemical sequence. Here, the average enthalpy of mixing is trivial: what matters is the fluctuations of composition inside one chain. This question has been analysed recently by A. D. Litmanovich and coworkers^[4] on a model with *two* chains in close contact: for this model, ref. ^[4] provides a detailed statistical analysis.

Our aim here is to extend the discussion to a real melt, in 3 dimensions, with many chains in interaction. This problem was solved long ago, at the level of the random phase approximation ^[5], by Frederikson, Milner and Leibler. ^[6] They set up a Landau theory, where

the order parameter is the deviation from average of the A/B concentration. For our case, with no correlations along the chemical sequence, they find a spinodal threshold when $\chi > 2$.

Our main aim, in the present note, is to present a poor man's approach to this segregation problem. This is crude, but possibly helpful, because it gives a certain insight for the resulting mesophase.

Weak segregation

We assume that the sequence A...BB...A of each chain is generated independently. (This is very different from a melt of polypeptides obtained from one same protein, where all the chains have the same sequence). The fraction ϕ of monomers A on one chain has a gaussian distribution around $\phi = 1/2$, with a r.m.s. width

$$\Delta = N^{-1/2} \quad (1)$$

We divide this distribution into two pieces:

(α) with $\phi > 1/2$

(β) with $\phi < 1/2$

The (r.m.s.) excess fraction of A in the (α) group is Δ . In the (β) group it is $-\Delta$.

We call $\psi(r)$ the local volume fraction of the (α) chains (ranging from 0 to 1). The average Flory Huggins energy ^[1] f of the $\alpha\beta$ mixture (per monomer site) is given by:

$$f/kT = N^{-1} [\psi \ln \psi + (1 - \psi) \ln (1 - \psi)] - \chi \Delta^2 (\psi - 1/2)^2 \quad (2)$$

This leads to a transition for:

$$\chi > 2N\Delta^2 \equiv 2 \quad (3)$$

Spinodals

Let us now focus on the case where χ is slightly larger than 2. We start at time 0, with a completely disordered melt. Then we expect to observe a spinodal decomposition ^[5] with a certain characteristic correlation length ξ_s . The formula for ξ_s may be obtained by adding to eq. (2) a term of order $a^2 (\nabla\psi)^2$ (where a is the size of one unit in the Flory picture). The coefficient of this gradient term is unaltered by our transformation from ϕ to ψ : in the language of ref. ^[5], the unperturbed correlation function for one chain is always the Debye function $g_D(r)$ ^[7].

The result (for χ slightly above 2) is:

$$\xi_s = a/3[2 - \chi]^{-1/2} \quad (4)$$

This gives the size of the regions which should nucleate initially. Indeed, a recent simulation of Houdayer and Muller ^[8] (with $N \sim 20$) shows A rich and B rich regions which may correspond to this.

At very long times, the α and β regions should grow into droplets, which swell by Ostwald ripening. But it is important to realise that the diffusion constants D should be slowed down by the weakness of the constant between A and B. We expect:

$$D \sim D_0 (\chi - 2) \Delta^2 \quad (5)$$

where D_0 is a diffusion constant for $\chi = 0$, which may be of the Rouse type (for small N) or of the reptation type (for large N).

Discussion

- 1) We see from eq. (3) that weak segregation should occur only for large values of the Flory parameter $\chi > 2$. This corresponds to very strong enthalpies of mixing. Thus all our discussion should be relevant only for some very special AB pairs, possibly using ionic components.
- 2) We should emphasize the weakness of the proposed segregation: the contrast between (α) and (β) regions is of order $\Delta \ll 1$, and the scattering intensities vary like $\Delta^2 \sim N^{-1}$.
- 3) We consider spatial variations in connection with eq. (4). But we can treat only wave vectors q which are small ($q R_0 < 1$, where $R_0 = N^{1/2} a$ is the size of one coil). At $q R_0 > 1$, the physically relevant value of Δ is $\Delta = g^{-1/2}$, where g is the number of monomers in a blob of size $q^{-1} = g = (qa)^{-2}$.

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